

The mathematical link between stratigraphic grain size variation of fall deposits and its time variation at the source

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Pyroclastic fall deposits which are produced by explosive volcanic eruption have various information on the eruption events. The areal distributions of pyroclastic fall deposits such as maximum grain size, median grain size, thickness, mass per unit area, etc. reflect the intensity in a single eruption and the wind conditions. Stratigraphic grain size variations of pyroclastic fall deposits also reflect the temporal behavior of the eruption intensity. For example, normal or reverse grading structures in the pyroclastic fall deposits have been attributed to temporal variation in the volcanic intensity (column height) and/or in the initial grain size distribution at the vent. However, no quantitative methodology has been developed to relate the temporal variation of source characteristics (column height and initial grain size distribution) to stratigraphic variation of grain size distribution at the deposits. In this study, we consider the mathematical description in 1D fall-sedimentation process, which relates the temporal variation of source grain size distribution to stratigraphic variation of grain size distribution.

The number of grains in a size bin must be conserved during sedimentation process and results in the same value at the arrival time on the deposits. The number of a specific-size grains between at the fallout time and at given times is linked by Lagrangian description. The key point is that the grain size and the departure time at the source are mutually related to the grain size and the arrival time at the deposits. As the arrival time corresponds to the stratigraphic location at the deposit, the stratigraphic variation of grain size in the deposit can be connected to the grain size characteristics and departure time at the source using the condition of grain number conservation. As a result, when the time variations of source grain size distribution and of fallout height are given, we obtain the temporal variation of grain size distribution at the sedimentation surface. It means that different sizes of grains which settle at the same arrival time are traced back to the different source time and height. The arrival time on the deposit can be related to the stratigraphic height in the deposit by the differential equation of increasing rate of the thickness, which equals to the volume flux through the sedimentation surface. By using these mathematical descriptions, we develop the mathematical method to link the temporal variation of eruption intensity to the stratigraphic variation of grain size distribution in the fall deposits.

Keywords: pyroclastic fall deposits, grain size distribution, stratigraphic variation, development of eruption

Color and grainsize of ash samples collected continuously at Sakurajima volcano, Japan

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Sakurajima volcano, Southwest Japan, is one of the most active volcano, and one of the leading volcano that are characterized by well-established geophysical observation network and enormous amount of data accumulation. Thus, the changes of magma plumbing system and the explosion processes at Sakurajima are now relatively well-understood for the time scale of years and of several hours, respectively (Iguchi et al., 2013). On the other hand, Sakurajima volcano has experienced several giant explosive eruptions every few hundred years, with some plinian eruptions and with erupted volume in the order of $>10^8$ m³. However, we do not understand the mechanism of phase transition from the recent small eruptions to such bigger ones. Even though we now reached a certain level to recognize that our knowledge of “vulcanian eruption” was too simple relative to real one, the fundamental processes that lead to such eruption variety are not yet constrained well enough. As the variety of eruption styles occurs in a time scale of months to days, we have been focusing on the sample acquisition in this time range and collected daily ash samples for years.

We have been succeeded in collecting daily samples by automatic sampling system for more than five years at the site ca 2 km south of Showa crater, Sakurajima volcano (Shimano et al., 2013). We also analyzed matrix glass composition to track chemical evolution of the magma system, and found a shift of FeO*/MgO in fall of 2009 before the waxing activity toward 2010. On the other hand, we have been searching for some real-time technique of petrological data for the comparison with geophysical data. The development of useful colorimeter and some results of heating experiments of ash resulted in understanding relationship between color of ash and condition of magma at depths (e.g., Yamanoi et al., 2008; Miyagi et al., 2013). So we have made time-series color measurements of ash samples for years. We found L* value, degree of brightness, decreased broadly during waxing stage in 2009-2010 whereas a* and b* values increased broadly at first several months but decrease gradually. These changes can be explained by increase of black fresh lava block particles, increase at the first stage and gradual fluctuating changes of old red oxidized particles in ash sample. We also measured color change of some grain size group and compared with the color of bulk samples. The preliminary results show that the color differs with particle size as Miyagi et al. (2010) reported, but a correlation was found between the color values of bulk sample and those calculated from the weight fraction and each color values.

Keywords: volcanic ash, continuous observation, color change, Sakurajima volcano

Numerical Simulation of Volcanic Ash Transport for the Eruptions at Mt. Shinmoe-dake during 26-27 January 2011

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The volcanic ash transport associated with the eruptions at Mt. Shinmoe-dake during 26 to 27 January 2011 is simulated using Japan Meteorological Agency Non-Hydrostatic Model (JMA-NHM) to verify the model with satellite observation. In the model, the mixing ratio and number concentration of ash particles are prognosed with the advection, diffusion, sedimentation, and source terms to represent the behavior of ash cloud. Simulation has been performed in the calculation domain covering 2500 km x 2000 km wide area with the horizontal resolution of 5 km.

The model is coupled with one-dimensional eruption column model to define the source term of ash particles, which is simply given as a function of the column height, the level of the release point, and the size of released particle, following Suzuki (1983) and Shimbori et al. (2010). Although the simulated distribution of ash cloud roughly agrees with satellite observation, close examination of the simulation result shows that the model fails to reproduce some of the ash clouds observed by the satellite, which means that much room still remains for improvement in the eruption column model in terms of release point and size spectra of ash particles. Three-dimensional direct numerical simulation has been conducted on a major sub-Plinian eruption during the period at Mt. Shinmoe-dake (Suzuki and Koyaguchi, 2013), in order to make new eruption column model with more realistic function for the source term of ash particles. As a result, it is found that the maximum release rate of the ash particles smaller than 100 μm appears in the height lower than that predicted by the usual eruption column model for same column top height. The authors are developing new eruption column model with realistic profile of release rate, based on this result, so as to improve the reproducibility of the ash transport with JMA-NHM. The sensitivity of the ash transport to altering the new and usual eruption column models will be presented.

Acknowledgement

This study was supported by the Earthquake Research Institute cooperative research program.

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Keywords: volcanic-ash dispersal, Atmospheric Transport Model,, Shinmoe-dake volcano, 2011

Weather Radar Investigation of Volcanic Smoke for Disaster-Prevention

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Sakurajima volcano has been active since February 2009. The total number of explosive eruptions was 966 in 2011, which was the highest number in recorded history. Corresponding to volcano activities, the ash accumulation in Kagoshima city increased and total ash amount of 3,500g/ m² was recorded at Kagoshima local meteorological observatory. Because the volume of volcanic ash in rural area paralyze public ground transportations such as rail road and highway, fast recovery efforts are required to the railroad company and city government. However, no quantitative volcanic ash fall estimation has been established. The present study focuses on utilization of operational weather radar for quantitative ash estimation (QAE), quantitative ash forecasting (QAF), and utilization of crustal movement information for providing ash volume which is necessary for initial conditions of a numerical diffusion model. Although the target volcano of the present study is Sakurajima, the knowledge on volcanic ash and algorithm developed by the present study can be applied to any other volcano which is located in operational weather radar observation area.

Keywords: weather radar, volcanic ash, quantitative ash estimation, Sakurajima, polarimetric radar

Eruption types determined by the mass flux and volatile component content of ascending magma flow

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Volcanic eruptions include the effusive ones that erupt fluid lava and the explosive ones that emit magma fragments in gassy flows or other forms. Traditionally, explosive eruptions are further classified into Plinian eruptions accompanied by high eruption columns, Pelean eruptions involving abundant pyroclastic flows, Vulcanian eruptions with strong instantaneous explosions, Strombolian eruptions involving periodic lava fountains and so on. Effusive and explosive eruptions are clearly controlled by the efficiency of degassing, but what produces various types of explosive eruptions has not yet been explained quite clearly. In this paper I propose a simple idea about how various explosive eruptions arise depending on the natures of ascending deep magma flow based on a stationary conduit flow model.

The state of erupting magma may be represented by exit velocity and the volatile content that determines the vesicularity of volcanic products. The volatile content is specified by its mass ratio to the fluid magma (including solidified part). On the other hand, it is convenient to represent the deep state of magma flow by the mass flux of fluid magma and the volatile contents before degassing takes place. When the magma flow is in a stationary state the mass flux of fluid magma is constant so that it defines the deep state of magma flow with the initial volatile content independently of the specific depth at which magma ascent starts.

The relation between the surface and deep states of ascending magma is calculated using a stationary conduit flow model. In this calculation the volatile component is assumed to move at the same speed as the fluid magma neglecting relative motions. In bubbly magma horizontal permeable flow of volatile gas is assumed to control the rate of degassing. In this treatment, the pressure gradient that drives the permeable flow is considered to arise from the ascent velocity change from center to side and the resulting difference of relaxation of gas expansion due to decompression (Ida, JVGR, 162, 172-184, 2007). The wall friction is assumed to be proportional to ascent velocity in bubbly flow and to the square of ascent velocity in gassy flow with suitable friction coefficients. The relation for water steam in magma is used for solubility of volatiles in magma.

The integration of conduit flow is executed from the surface to a deep conduit. Namely, the deep state of magma flux and volatile content are calculated for various sets of the exit velocity and volatile content at the surface prescribed with the magma pressure equal to the atmospheric pressure. Compiling the calculation results shows that some groups characterize the relation between the surface and deep conditions. Each group can be interpreted in connection with eruption types in the following way.

Firstly, high-speed gassy flow erupts violently when the deep magma contains sufficiently abundant volatile component. This case may produce a Plinian eruption. In this case the exit velocity and gas content are determined by the deep magma flux alone independently of deep gas content because of adjustment by degassing during the ascent process. Secondly, a stationary conduit flow disappears below the critical value of fluid magma flux with high gas content in a deep conduit. In this case magma flow should be non-stationary and may produce Vulcanian or Strombolian eruptions. Thirdly, a relatively slow magma flow with low vesicularity flows out when volatile component is poor. This case likely results in Pelean eruptions because of difficult acceleration of gassy flow in the air.

Our analysis and interpretation suggest that various eruption types arise from different combinations of ascending magma flux and degassing efficiency. It is non-linearity involved in ascending magma flow with vesiculation and degassing that produces separate groups characterizing eruption types.

Keywords: volcanic eruption type, ascending magma flow, conduit flow model, volatile component content, degassing, computer simulation

Magma eruption rates, eruption styles, and preeruptive magma viscosity

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Magma eruption rate is one of the most fundamental parameters for a volcanic eruption (e.g., Pyle, 2000). It is obtained mainly by geophysical or geological observations. We interpret this important parameter from a petrological point of view and also from a fluid dynamic point of view.

We have collected a hundred of data on magma eruption rates, bulk rock chemical compositions, and phenocryst contents for various styles of eruptions (Plinian, sub-Plinian, basaltic Plinian, lava flow, and lava dome). We are compiling these data on the basis of their 'preeruptive magma viscosities', which are important measures of magma eruptibility (Takeuchi, 2011). Preeruptive magma viscosity is the viscosity of magma (melt + crystals) in the magma chamber at the preeruptive conditions. This value can be obtained by the bulk rock chemical composition and phenocryst content, using an empirical formula (Takeuchi, 2010). We have found that eruption styles are closely correlated to preeruptive magma viscosities but poorly correlated to bulk rock compositions.

We have also examined the difference in magma eruption rates between the explosive phase(s) (e.g., Plinian) and the effusive phase (e.g., dome) in a series of eruptions, in order to understand the transition between these two eruption styles (e.g., Kozono and Koyaguchi, 2009a,b). We have found that the difference is positively correlated to preeruptive magma viscosity.

The above results indicate that preeruptive magma viscosities largely control eruption styles and eruption rates. Our results also show that the eruptive magmas are divided into two types, low-viscosity type (basalt to low-phenocryst-content andesite) and high-viscosity type (high-phenocryst-content andesite to rhyolite). The boundary is at about 10^4 Pa s. These two types may be closely linked to the magma generation processes (e.g., fractional crystallization and melt segregation from crystal mush).

Keywords: magma eruption rates, eruption styles, preeruptive magma viscosity, transition between explosive and effusive eruption, phenocryst content, bulk rock chemical composition

Conditions for transition from lava dome to explosive eruption

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Conduit flow dynamics involving magma vesiculation, gas escape, and crystallization during a lava dome eruption lead to complex processes such as a transition to an explosive eruption. Because the transition from the lava dome to the explosive eruption is accompanied with a drastic increase in eruption intensity, it is important for volcanic hazard mitigation to determine conditions for this transition to occur. In this study, on the basis of a 1-dimensional conduit flow model, we investigated how the conditions for the transition from the lava dome to the explosive eruption depend on magmatic and geological parameters.

In order to systematically investigate the dependence of the transition conditions on the magmatic and geological parameters, we used the relationship between chamber pressure (p_{ch}) and mass flow rate (q) for steady conduit flow (the p_{ch} - q relationship). When the slope of the p_{ch} - q relationship (dp_{ch}/dq) has a positive value (positive differential resistance), the steady flow is stable. When dp_{ch}/dq has a negative value (negative differential resistance), on the other hand, the steady flow is unstable. The negative differential resistance is generated by two positive-feedback mechanisms. First, effective magma viscosity decreases with increasing q because of delay of crystallization, leading to reduced viscous wall friction (feedback 1). Second, magma porosity increases with increasing q because of less efficient gas escape, leading to reduced gravitational load (feedback 2). These two feedback mechanisms induce a sigmoidal p_{ch} - q relationship for some realistic conditions; the positive differential resistance in the low- q and high- q regimes, and the negative differential resistance in the intermediate regime. The analyses of time-dependent conduit flow model indicate that, because of the sigmoidal p_{ch} - q relationship, as magma supply at depth gradually increases from the low- q regime to the intermediate regime, magma discharge rate abruptly increases from the low- q to high- q regimes. This abrupt increase in magma discharge rate accounts for the transition from a stable lava-dome eruption to an explosive eruption. We, therefore, define the value of q at the boundary between the low- q and the intermediate regimes as the critical magma supply rate for the transition (q_{cr}).

Our results show that q_{cr} is mainly controlled by the feedback 2 for a wide range of magmatic and geological conditions, whereas it is controlled by the feedback 1 only when phenocryst content is very high. When q_{cr} is controlled by the feedback 2, the value of q_{cr} depends on parameters related to gas escape such as the permeability for vertical gas escape and that for lateral gas escape. We found that for a plausible range of vertical permeability which is constrained from permeability measurements of volcanic rocks, q_{cr} remarkably decreases with decreasing lateral permeability, and it becomes substantially lower than typical magma discharge rates for observed lava-dome eruptions in the limiting case of zero-lateral permeability (i.e. no lateral gas escape). This indicates that the presence of lateral gas escape is a necessary condition for a stable lava-dome eruption to occur. In addition, we found that q_{cr} strongly depends on conduit radius owing to the effects of the change in the conduit radius on the degree of gas escape. As the conduit radius decreases, the ascent of the liquid is suppressed because of the increase in wall friction force, which promotes vertical gas escape. The decrease in the conduit radius also induces an increase in the ratio of the perimeter to the cross-sectional area of the conduit and a decrease in the length scale of pressure gradient that drives lateral permeable gas flow, which promotes lateral gas escape. These promotions of gas escape lead to an increase in q_{cr} . The above results suggest that the variation of conduit radius is a key factor for the transition from a lava-dome to an explosive eruption.

Keywords: conduit flow, numerical model, eruption transition, lava dome, explosive eruption, gas escape

A phreatic explosion model for Mayon volcano, Philippines, inferred from analyses of an explosion earthquake

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Mayon is one of the most active volcanoes in the Philippines with 49 known historical eruptions from 1616 to 2010. A phreatic explosion took place at Mayon on 7 May 2013 that killed five climbers. In this presentation, we show the results of the waveform inversion for the explosion earthquake and discuss a phreatic explosion model for Mayon.

During the explosion in 2013, a VLP event with a peak frequency of 0.4 Hz was recorded by three broadband seismometers which we had installed in 2011. We performed a frequency-domain waveform inversion in 0.1-0.6 Hz, which pointed to a combination of a subhorizontal tensile crack and a vertical single force at a shallow part beneath the summit crater. Contributions from the crack and force to the waveforms had amplitudes comparable to each other.

The source time functions obtained by the waveform inversion are bandpassed forms (filtered source time functions; FSTFs), which may be different from the source time functions without filters (deconvolved forms of the source time functions; DSTFs). Instead of performing numerically unstable deconvolution operations, we assumed simple step- and impulse-type functions with finite durations as candidates of the DSTFs. We applied the bandpass filter to these functions and compared with the FSTFs. The bandpassed waveforms of the impulse-type functions were similar to the FSTFs for both the crack and force, suggesting that the DSTFs can be approximated by the impulse-type functions. The estimated DSTF for the tensile crack showed an inflation followed by a deflation, whereas that for the single force showed a downward force around the time of the maximum opening of the crack.

The RMS seismic amplitudes, GPS baseline lengths over the volcano, ground surface temperatures around the summit, waveform correlations among the seismic stations, the sulfur dioxide emission, and rainfall did not show clear precursory signals.

Since the analyzed VLP event occurred during the phreatic explosion, the initial inflation in the DSTF for the crack may have been caused by boiling of underground water. This crack is subhorizontal and located at a shallow part, suggesting that the crack may be located on a boundary between permeable and impermeable layers where the water may have accumulated and finally boiled, generating the explosion. The downward force may represent the counter force of the explosion. The deflation of the crack in the latter half of the DSTF may have been caused by outgassing of water vapor during the explosion.

The estimated moment amplitude is explained by a volume change of $400 \text{ m} \times 400 \text{ m} \times 0.4 \text{ m}$. A topographic change comparable to this crack size was not observed during the explosion, suggesting that the explosion destroyed only a limited portion of the crack. This crack may repeat the explosion once the fragmented portion of the crack is sealed through the hydrothermal alteration. At Mayon, small ash explosions occurred in 2003, 2004, 2006, and 2009, with intervals of a few days to a few years (mostly longer than one month). These intervals are close to experimentally derived time scales of fracture sealing by the hydrothermal alteration; according to the experiments by Berger et al. (1994, *Geochim. Cosmochim. Acta*), who used basalt samples with chemical compositions similar to magmatic eruption deposits at Mayon, a centimeter-order fracture is sealed in two months for 300 °C and 54 months for 150 °C.

In this model, timings of the explosions are controlled by the sealing of the fracture. Therefore the explosions can occur even with a constant supply of heat and water. This may be the reason why no clear precursory signals were observed before the explosion in 2013. No VLP events other than that associated with the explosion in 2013 have been observed at Mayon since the beginning of the observation in 2011, which may be explained by the sealing of the fracture just prior to the explosion.

Keywords: Phreatic explosion, Waveform inversion, Source time function, Hydrothermal alteration

Magma accumulation process of Izu-Oshima volcano, as revealed from deep LF earthquakes, deformation and CO₂ out-gassing

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In order to make successful mid-term eruption predictions, we need to detect particular precursory processes operating in magma-plumbing system. Since 1989, Izu-Oshima volcano has continued its re-inflation, after the last eruption in 1986, and further repeated deflation-inflation cycles, resulting a net inflation of the volcano. The rate of secular inflation decreased exponentially until 2006, while the amplitudes of the deflation-inflation cycles increased. Since 2007, the rate of secular inflation has kept a constant speed and has also increased the activity of deep low-frequency (LF) earthquakes occurring at the depth range of 30-40 km beneath the volcano. Each episodic LF earthquake activity was preceded by the volcano deflation and accompanied by the inflation. Based on these evidences, we may suppose that the volcano inflation is caused by the supply of magma from a source region at the depth range of 30-40 km beneath the volcano, and that an episodic out-gassing from the shallow magma reservoir triggers each deflation-inflation cycle. To demonstrate the proposed mechanism, we need to combine the data on magma accumulation and out-gassing processes. To monitor the out-gassing of basaltic magma accumulating beneath the volcano, CO₂ is most helpful. In September 2005, we started continuous monitoring of soil CO₂ concentration at the summit of Izu-Oshima volcano, and obtained an evidence for the out-gassing process; the correlated increase of soil CO₂ concentration during the periods of not only accelerated inflation but also deflation of the volcano. Integrating the observational data, we suppose that the rate of magma supply from the upper mantle has increased since 2007 and that the increase in amplitude of deflation-inflation cycles might indicate a volume increase of CO₂ over-saturated region at the upper part of the magma reservoir beneath the volcano.

Keywords: Izu Oshima volcano, eruption prediction, precursors to eruption, volcano deformation, CO₂ out-gassing

Spatio-temporal characteristics of volcanic tremor during the 2011 Kirishima eruption by seismic wave analysis

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Volcanic tremors are considered as oscillations occurred in magma supply system and provide us important information about condition of the system from magma chamber to crater through a conduit. Therefore, it is important to investigate locations of their source and their characteristics for understanding the condition and process of volcanic activity and modeling the magma supply system. In this study, we reveal features of volcanic tremor using seismic data at Kirishima volcano.

An array observation help us to get information of incident waves on the stations. The two array seismic observations were carried out around Shinmoe-dake during the 2011 Kirishima eruption. One consisted of 25 seismometers located 3 km southwest of the Shinmoe-dake crater (Matsumoto et al., 2013), and the other consisted of 16 seismometers located 5 km northeast of the crater (Nakamichi et al., 2013). The combining data from two arrays enable us to determine the tremor sources. Moreover, we estimated mechanism of the tremor source by using waveforms recorded at temporal seismic station in Kirishima volcanic area and tremor sources obtained by array analysis.

This study focus on the volcanic tremor which occurred on February 2, 2011, and its duration was about 40 minutes. Peak frequencies of the tremor were about 1, 2, 3, and 4 Hz. We investigated temporal variation in the source location of the tremor from the slowness and azimuth of incident wave by MUSIC method. we found that most part of the tremor were radiated around Shinmoe-dake crater. In this part, the tremor that had large slowness and relatively long duration was located in shallow region beneath the crater. In contrast, at some parts of the tremor, source location for waves with short duration were near Ohnami pond, 3.3 km northeast of the crater. Based on amplitude analysis for the seismogram recorded by the seismic network, we also found out difference in radiation patterns of the volcanic tremor among the tremor sources. Assuming single crack model, we found the strike and dip direction of the crack beneath the Shinmoe-dake crater is different from those near Ohnami pond.

Keywords: volcanic tremor, Shinmoedake

Relationship BH-type earthquake swarms and ground deformations prior to eruptions at Showa crater at Sakurajima volcano

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Vulcanian eruptions have occurred at the Minamidake crater at the summit since 1955 at Sakurajima volcano. Principal eruptive activity shifted to the Showa crater at the eastern flank of the summit in 2006. The eruptions at the crater become active and minor vulcanian eruptions occurred about 1,000 times per year in 2010-2013. Ash plume height of the eruptions sometimes reached to 3000m in 2013. The eruptive activities at the Showa crater gradually increase.

Inflationary strain changes are observed by extensometers a few tens of minutes to several hours prior to the eruptions and are caused by pressure sources located at depths of 0-1.5 km (Iguchi et al., 2013). The inflation rates decrease or sometimes suspend about 30 minutes before the eruptions. Small earthquakes dominated by high frequency components (5-6 Hz) swarm when duration of inflation is longer than 1 hour. The hypocenters of the earthquakes are located at a depth of 0.5 km beneath the crater and are close to depth of the pressure source. The earthquakes begin to occur a half hour to 1 hour after the start of the inflation. The amplitudes and number of the earthquakes further increase when the inflation rates decrease or suspend. And, the occurrences of the earthquakes suddenly stop at the start of the eruptions. The occurrences of the earthquake swarms are related to the decrease of inflation rate and the long inflation.

Seismic energy releases of the precursory earthquakes related to every eruption accelerate before eruptions. There are all kinds of large and small seismic energy releases in the eruptions. The accelerations of the seismic energy releases before eruptions with explosive events tend to be larger than those with non-explosive events. And, the accelerations of the seismic energy releases are rapid in the case of large deflations after eruptions. The precursory earthquakes may be generated by release of excess pressure accumulated by inflation of the pressure source. We might be able to predict eruption types and scales from occurrence patterns of the precursory earthquakes.

Keywords: Sakurajima volcano, precursory earthquake, ground deformation

Volcanic tremor caused by flow-induced oscillation of a magma-filled dike

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Volcanic tremor (VT) is known to be long-period and long-duration ground motion generally observed during or before eruptions. Majority of VTs show an emergent onset and, when accompanied with eruptions, an exponentially growing phase in amplitude is typically observed (Konstantinou and Schlindwein 2002; McNutt and Nishimura 2008). This characteristic suggests that VT is manifestation of self-oscillation, in which a persistent steady forcing excites an eigen oscillation of a system and the amplitude exponentially grows until a nonlinear process leads to a limit cycle. In volcanic settings, a steady magma flow through an underground conduit may cause flow-induced self-oscillation of bedrocks, and this is the idea first presented by Julian (1994). In the case of the collapse of Tacoma Narrows bridge, which is known to be caused by flow-induced oscillation, the vibrating bridge can be modeled by an elastic plate placed parallel to an infinite flow. Here we consider a reversed setting: a fluid-flowing thin layer in an infinite elastic body. This system is also unstable if the flow speed is high enough, and may be a generation mechanism of some VTs.

We consider a fluid-flowing plane layer sandwiched between two semi-infinite elastic bodies, expanding Julian's idea to a more general elastodynamic model (for details, see Sakuraba and Yamauchi 2014 to appear in *Earth, Planets and Space*). The eigen oscillation that should be excited in this self-oscillation model is an elastic surface wave. Therefore, we solve the Navier-Stokes equation linearized about a laminar flow with the boundary condition that can maintain a surface wave traveling along the layer. We succeeded in obtaining a complex phase speed of the surface wave as a function of wavenumber (and some parameters) using a shooting method, and found that a relatively slow magma flow could lead to instability in which the imaginary part of the phase speed is positive. Remarkably, the most unstable mode exhibits an antisymmetric (flexural) deformation, which has not been discussed in previous similar studies (Balmforth, Craster and Rust 2005; Dunham and Ogden 2012). The unstable mode is identical to two parallel Rayleigh waves traveling against the basic magma flow. The instability can be understood as acceleration of nearly circular particle motion of the Rayleigh wave due to viscous drag of the main laminar flow. As the critical flow speed giving a neutrally stable state decreases in inversely proportional to wavelength, instability will occur with the largest possible wavelength, which could be several kilometers and produce an oscillation period of around 1 second. In that case, the critical flow speed can be reduced to less than 1 m/s when the magma is basaltic. As the magma flow speed in a dike will not exceed several meters per second, there would be a lower bound in the critical wavelength, producing a period of 0.1 second. Thus our model naturally explains typical periods of 0.1-1 seconds observed on volcanoes. Our model also explains typical timescales of the linearly growing phase in some VTs.

Magnitude-frequency distribution of volcanic eruptions from an open conduit

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Explosive eruptions such as vulcanian, strombolian, or significant gas burst excite seismic waves. Such seismic signals, which are often called as explosion earthquakes or explosion quake, are used to quantitatively evaluate the magnitude of such volcanic explosions. In the present study, we systematically examine magnitude-frequency distributions of explosion earthquakes observed at Sakurajima and Suwanosejima volcanoes in Japan, and Semeru and Lokon volcanoes in Indonesia. We use the long-term catalog data of Sakurajima explosions for the period from 1963 to 2011, which are routinely summarized by Sakurajima Volcano Research Center (SVRC). Also, we measure the amplitudes of explosion earthquakes from continuous seismic records observed at Suwanosejima and Semeru and Lokon volcanoes in Indonesia. We measure the number of earthquakes that exceed a given amplitude, and then plot the cumulative number of earthquakes versus amplitude, as is often done for examining Ishimoto-Iida's relation, which expresses a power law distribution. However, the observed frequencies of earthquakes at the four volcanoes do not seem to fit the Ishimoto-Iida's relation. The cumulative numbers are well explained by exponential functions. This means that the magnitude of explosion earthquake at each volcano is randomly determined with an average scale.

Keywords: Volcanic Explosion, Vulcano, Gas burst, Magnitude-frequency distribution

Rheological transition of plagioclase-bearing magma: high-temperature uniaxial deformation experiments of sanukite lava

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High-temperature uniaxial compression experiments were done for bubble- and phenocryst-free, plagioclase-bearing lava to investigate the effect of tabular crystals on rheological properties of highly crystalline magma. High-Mg andesite lava from Goshikidai lava plateau, southwest Japan, was used for starting material. The lava is bubble- and phenocryst-free, composed of 60 vol.% rhyolitic glass, 36 vol.% of tabular plagioclase and 4 vol.% of pyroxenes and magnetite and plagioclase crystals are well aligned parallel. High-temperature uniaxial deformation apparatus at Earthquake Research Institute, the University of Tokyo, was used for experiments. The lava was cut to 10 x 10 x 20 mm rectangular solids and deformed under conditions of temperatures of 1238, 1188, and 1138 K and deformation rates from 3.16 to 0.003 mm/min. Run samples were quenched and processed to thin section for textural and compositional analyses by using EPMA.

Phase proportions in all run samples were the same as that of starting material, indicating crystallization did not occur during experiments. The lava behaves as shear thinning fluid under all temperature conditions. Viscosity at strain rate of 10^{-4} s^{-1} increases from $10^{8.7}$ to $10^{9.4}$ Pas with decreasing temperature. Power law exponent [= $d(\log \text{ viscosity})/d(\log \text{ strain rate})$] is ca. 0.64, which is consistent with extrapolation of previous studies for natural plagioclase-bearing magmas. Relative viscosity [= (bulk viscosity)/(melt viscosity)] is ca. $10^{2.4}$ at strain rate of 10^{-4} s^{-1} under all temperatures, indicating that the concept of relative viscosity works well under present experimental condition. The relative viscosity-crystal fraction relation is also consistent with extrapolation of previous studies for natural plagioclase-bearing magmas. Marron-Piece equation well explains the relation with the maximum packing fraction of 0.43. Present results suggest that rheological transition occurs at crystal fraction near 0.43 for plagioclase-bearing natural lava in which plagioclase crystals are well aligned parallel. The value is higher than ca. 0.3 proposed by Picard et al. (2013)'s experiments in which plagioclase orientation is random in starting materials, indicating the first order importance of plagioclase orientation distribution on rheological transition.

Keywords: rheological transition, viscosity, magma, plagioclase, non-Newtonian fluid, sanukite

Conduit flow of silicic magma: Viscous flow or Frictional sliding?

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Outgassing rate and bulk magma viscosity that control the style of volcanic eruptions depend on flow type of magma ascending in a volcanic conduit. When magma behaves as a Newtonian fluid, magma in the conduit experiences shear strain large enough to cause effective outgassing. On the other hand, once shear starts to localize, bulk magma viscosity may decrease due to slip deformation and outgassing rate also decreases in parts other than shear-localized region (Okumura et al., 2013 EPSL). Silicic magma experiences shear-induced brittle fracturing and subsequent frictional sliding along the fracturing zone during its ascent (e.g. Gonnermann and Manga, 2003 Nature; Tuffen et al., 2003 Geology). Therefore, outgassing rate and bulk magma viscosity during the ascent are expected to change dramatically. Previous studies (Tuffen et al., 2003 Geology; Gonnermann and Manga, 2005 EPSL) also proposed that fractured magma can heal during magma ascent and that fracturing and healing processes may control the dynamics of magma ascent. In contrast to this model, some experimental studies (e.g. Okumura et al., 2010) indicated that fractured magma cannot heal as long as the deformation continues. In this study, we perform deformation experiments for fractured magma to investigate flow type of magma in the conduit, i.e. viscous flow or frictional sliding, and controlling factors of the transition from viscous flow to frictional sliding.

The deformation experiments were carried out using a custom-made torsional deformation apparatus which was installed in synchrotron radiation X-ray imaging system (BL20B2) of SPring-8. To simulate fractured silicic magma, we crushed rhyolite obsidian and sorted them into fragments of 75-250 μm in size. The powdered sample was sandwiched by two obsidian discs and they were twisted by rotating a piston attached with a rotational motor. The torsional deformation experiments were performed at temperatures of 800 and 900 $^{\circ}\text{C}$ under 1-10 MPa pressures. The rotational rate was set to be 0.1 to 10 rpm, corresponding to strain rates of 10^{-2} to 1 s^{-1} if the sample deforms homogeneously. The deformed samples were observed in situ using an X-ray radiography.

At a temperature of 900 $^{\circ}\text{C}$ and rotational rates of 0.1-1 rpm, homogeneous deformation through a sample was observed under a pressure of 10 MPa, which indicates viscous deformation. In contrast, the sliding at the interface between powdered obsidian and the disc was observed under 1 and 5 MPa pressures. At a temperature of 800 $^{\circ}\text{C}$, the sliding was found under 1-10 MPa pressures. These results indicate that frictional sliding along fractured zone is flow type of magma in shallow parts of the conduit (<10 MPa).

We assume that flow type is determined by competition of shear stress necessary for viscous flow and frictional sliding. If magma has high viscosity and shear stress to deform a sample viscously is large, the flow type becomes frictional sliding. At a temperature of 900 $^{\circ}\text{C}$, viscous flow and frictional sliding were found at 10 and 1-5 MPa pressures, respectively. At this condition, magma viscosity is approximately 10^7 Pa s (Hess and Dingwell., 1996) and shear stress necessary for viscous deformation is 1 MPa at a strain rate of 0.1 s^{-1} . Because the frictional sliding was observed at pressures of 1-5 MPa, the frictional coefficient is estimated to be ca. 0.1. When we use this value and the criterion for shear-induced brittle fracturing proposed by Okumura et al. (2010), the dynamics of magma ascent is controlled by frictional sliding at shallow parts of the conduit. In addition, silicic magma can ascend quickly due to low frictional coefficient of fractured magma.

Keywords: Silicic magma, Volcanic eruption, Viscous flow, Frictional sliding, Synchrotron radiation X-ray

Mechanism of delayed brittle-like fragmentation of vesicular magma analogue

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Magma fragmentation is a key phenomenon controlling volcanic eruptions. The fragmentation is classified into two styles: solid-like brittle fragmentation and liquid-like ductile fragmentation. Brittle fragmentation is more hazardous than ductile fragmentation because violent release of pressurized gas in the bubbles contained in the magma may lead to explosive eruptions

The fragmentation of magma, which is a viscoelastic fluid, occurs through a combination of viscoelasticity and rapid deformation. We conducted a rapid decompression experiment over ten years in order to clarify the viscoelastic effect on the fragmentation using a magma analogue, syrup containing gas bubbles (Kameda et al GRL 2008; Kameda et al. JVGR 2013).

Through the experiment we demonstrated the existence of a transitional fragmentation behavior. This transition behavior, which we refer to herein as brittle-like fragmentation, occurred even if the response of material should be in a ductile manner. Comparing the realistic decompression time with the viscoelastic relaxation time for magma, it is more probable that the fragmentation in the real volcanic system occurs in a brittle-like manner.

Observation by high-speed photography using a visible light-source indicated that the onset of brittle-like fragmentation was triggered by the sudden release of a considerable amount of gas from a crack in the specimen. Further observation (Shida et al. IAVCEI 2013) showed that reducing the volume of the specimen suppressed the onset of fragmentation even if their brittleness (Ichihara and Rubin 2010) was close to unity. In our case, the pore distribution of the small samples was more uniform than that of large samples. This observation implies that the crack is initiated from the interior of the specimen due to non-uniform spatial distribution of bubbles.

Then, we observed the interior of the specimen by synchrotron X-ray tomographic microscopy. The X-ray tomographic microscopy was performed at the BL20B2 beamline of the Japan Synchrotron Radiation Research Institute (JASRI, Hyogo, Japan). Initial structure of the specimen was observed by three-dimensional tomographic imaging. High-speed radiography was performed during the decompression. A digital charge-coupled device (CCD) camera was used as the detector whose imaging area is about 16 mm (horizontal) by 5 mm (height) with spatial resolution of 8 $\mu\text{m}/\text{pixel}$. We took 1800 projections over 180 degrees of rotation for tomographic imaging. The framing rate of radiography is 200 frames per second. We successfully captured a series of images during the brittle-like fragmentation. The reconstructed 3D image of the specimen indicated that the brittle-like fracture was initiated at a notch and a chain of small bubbles in the vicinity of a large bubble.

We propose the following scenario for brittle-like fragmentation: It is initiated from ductile growth of internal cracks by connection of bubbles. The stress concentration and the brittleness at crack tip may exceed the critical level at some moment, which leads to brittle failure of the crack. Rapid decompression due to sudden release of a considerable amount of gas from the crack may increase local brittleness to cause partial fragmentation. Such a sequence from ductile crack growth to partial fragmentation may successively occur in brittle-like fragmentation.

Keywords: Magma, Fragmentation, Viscoelasticity, Decompression, X-ray CT

Approach by volcanic observation for dynamics of volcanic phenomena

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Dynamics of volcanic activity has been revealed by various kinds of observation. Inflation of volcanic body prior to eruptions corresponds to intrusion of magma to the underground of volcanoes and deflation is accompanied by eruptive activity.

Prior to vulcanian eruptions, minor deflation is detected. This corresponds to leakage of volcanic gas from a gas pocket formed at the uppermost conduit. The minor deflation reflecting minor pressure decrease induces a sudden degassing in oversaturated magma. The sudden degassing corresponds to a volume increase at a deep part in the conduit, as revealed by outward first motions of explosion earthquakes. The volume increase attained at the top of the conduit and the gas pocket overbursts. As the results of collapse of the gas pocket, infrasound of air-shock type is generated with ejection of incandescent bombs followed by ejection of volcanic ash.

Keywords: vulcanian eruption, Sakurajima

Tephra-Fall Predictions with the JMA Regional Atmospheric Transport Model for the 1914 Eruption at Sakurajima Volcano

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A hundred years ago, the largest eruption in 20th century Japan, *i.e.* the Sakurajima Taisho Eruption occurred on 12 January 1914. With this eruption, tephra-falls were observed large-area in Japan from Kyushu to Tohoku region (Hasegawa, 1914; Omori, 1916). For such large-scale eruptions, the Meteorological Research Institute (MRI) is going to do new research project of the tephra-fall predictions with the JMA Regional Atmospheric Transport Model (JMA-RATM) from this year. The research will lead to the improvement of the Volcanic Ash Fall Forecasts (VAFFs) of the Japan Meteorological Agency (JMA). For the purpose of investigating the predictability of the current JMA-RATM against large-scale eruptions, predictions of volcanic-ash dispersion and tephra-fall for the Sakurajima Taisho Eruption were carried out. The initial values which are the total volume (ash and pumice) of $6 \times 10^8 \text{ m}^3$, the plume height of 3000-18000 m and the eruption duration of 38 hours are assumed based on Yamashina (1999), Yasui *et al.* (2006), Iguchi (2014) and so on. The input GPVs are the JMA Mesoscale Analysis after 28 March 2013. The forecast time by the JMA-RATM is 72 hours from starting time. Results of the calculations indicate that, under the assumption of the ash-density of 1 g/cm^3 , the predictions of tephra-fall depths are over 1 m in Sakurajima Island for weak-wind weather condition, several 10 cm at Kagoshima City for easterly wind in summertime, and of the order of 0.1-1 mm at Osaka, Nagoya and Tokyo Metropolitan area for southwesterly wind field. In the atmosphere, dispersions of volcanic-ash up to Tohoku and Hokkaido region are also predicted at the same forecast time.

In the presentation, from the results of volcanic-ash dispersion and tephra-fall prediction, the predictability and problems of the current JMA-RATM for large-scale eruptions will be reported.

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Keywords: Atmospheric Transport Model, volcanic-ash dispersal, tephra fall, Volcanic Ash Fall Forecast, Sakurajima volcano, 1914

Evaluation of wind data for tephra dispersion simulations

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Understanding how pyroclasts disperse from volcanic plumes is a fundamental problem of volcanology to reconstruct eruption conditions from tephra fallout deposits. Tephra dispersion is not only a scientifically interesting but also socially and economically important problem. For this reason, advection-diffusion models for tephra transportation have been developed with simplified assumptions (e.g. TEPHA2, PUFF, FALL3D). Simulation results of these advection-diffusion models are affected by the input wind data. For example, in a case study of the Kirishima 2011 eruption with PUFF, simulation results with finer temporal-spatial resolution wind data (Japan Meteorological Agency Mesoscale Model and ERA Interim) reproduced a more wavy shape of observed eastward extending plume (about 900 km from the vent) than that with coarser temporal-spatial resolution wind data (NCEP/NCAR Reanalysis).

Some wind data are available from Japan Meteorological Agency Mesoscale Model, ERA Interim of the European Center for Medium-Range Weather Forecasts and NCEP/NCAR Reanalysis of the National Oceanic and Atmospheric Administration of the United States. Between these wind data, there are differences in data assimilation methods, forecast models and temporal-spatial resolution. In addition, a finer temporal-spatial resolution wind data can be generated with using weather forecast models, such as the Weather Research and Forecasting (WRF) model. The WRF model is a fully compressible, Euler non-hydrostatic mesoscale forecast model developed by a multiagency collaboration. The WRF model is suitable for use in a broad spectrum of applications across scales ranging from meters to thousands of kilometers.

We are developing a system to generate wind data suitable for simulations of the advection-diffusion models. Such wind data should have a spatial resolution of several hundred meters near the vent, that of several kilometers far from the vent and vertically several tens of layers. In addition, it must be required to reproduce the interaction between ambient atmosphere and volcanic plumes. For this purpose, we are carrying out numerical calculations with the WRF model and the available wind data sets; we attempt to generate wind data with higher temporal-spatial resolution using data assimilation based on the observations from the regions of interest (e.g. volcano locations and downwind area) and other observations (e.g. the radar observations).

Keywords: Tephra dispersion simulation, Weather forecast model

Classification of infrasound waveforms and analysis of video images at volcano eruption

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Volcano infrasound is an increasingly useful technique for measuring and analyzing eruptive activity. In order to reveal causes of characteristics of infrasound waveform, we analyzed infrasound records and video images at eruptions that occurred at Yasur Volcano, Vanuatu during an hour at 15:00, 21th September, 2011. Yasur Volcano has three active vents in the summit crater. Infrasound waves are generated every 1-3 minutes at various eruptions with reddish magma, vapor rich gas and ash. Infrasonic waveforms are mostly symmetric with a sharp compressive onset, followed by a small rarefaction phase. In this study, we analyzed infrasound events whose maximum amplitude is more than 50 Pa at one station on the summit crater rim. We conducted cross correlation analysis to 29 wave records and classified them into 3 types; A (13 events), B (12 events) and C (4 events). Time window of the analysis was 5s from a second before the maximum peak time. In contrast to type C, cross correlation coefficient among wave forms of types A and B was very high ($> \sim 0.75$). The compressive ratio and the ratio of positive and negative peak are different between type A and B. On the other hand, we analyzed selected video sequences of these eruptions. We read the RGB and brightness values on horizontal line above each vent until the end of the eruptions from 2s before the ejecta reaches the line. The maximum R and brightness value of type B events was twice as other type eruptions. This result may suggest that type B eruptions release relative much reddish magma and whitish gas compared to those of type A and C.

Keywords: infrasound, volcano eruption

Numerical treatment of dry bed problem in the model of pyroclastic flows based on the 1-D shallow-water equations

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During explosive volcanic eruptions, a mixture of pyroclasts and volcanic gases is released from the vent. When the mixture loses its upward momentum before the density of the mixture becomes lower than the atmospheric density, the mixture forms a pyroclastic flow. Dynamics of pyroclastic flows can be approximated by that of an inviscid gravity current. The dynamics of inviscid gravity currents are controlled by an inertial-buoyancy balance on the front (e.g., Benjamin, 1968); we refer to this condition as "the front condition". The front condition, and hence, the dynamics of the inviscid gravity currents strongly depends on the density ratio of the current (ρ_c) to the ambient (ρ_a) (e.g., Ungarish, 2009). When $\rho_c/\rho_a \sim 1$, the current is characterized by a high front, whereas a front height does not develop when $\rho_c/\rho_a \gg 1$. In pyroclastic flows, because density ratio ρ_c/ρ_a varies spatially and temporally, the dynamics of pyroclastic flows becomes complicated; the basic features of pyroclastic flows, such as the run-out distance, have not been fully understood. The aim of our study is to develop a unified model of the inviscid gravity currents for various density ratio ρ_c/ρ_a .

In general, the dynamics of shallow inviscid gravity currents can be described by the shallow-water equations. There are two numerical models to solve the shallow-water equations: "shock front condition model" (SFC model) and "artificial bed-wetting model" (ABW model). SFC model is a model, in which the front condition is applied to the boundary condition (e.g., Ungarish, 2009). The boundary condition is given as a function of ρ_c/ρ_a . On the other hand, in ABW model, an artificial wet bed with the height of ϵh_0 is set on the dry bed in order to express the front condition, where h_0 is a characteristic height scale (e.g., Toro, 2001; Larrieu et al., 2006; Doyle et al., 2007). This model has the only parameter ϵ for the front condition. Although the front condition, and hence the appropriate value of ϵ must be a function of ρ_c/ρ_a , the relationship between ϵ and ρ_c/ρ_a has not been studied. In order to resolve these problems, we carried out parameter studies using the two models for solving a simple "one-dimensional (1-D) dam-break problem".

On the basis of systematic comparisons between the results of SFC model and ABW model, we found the relationship between the parameter ϵ and ρ_c/ρ_a : $\epsilon \sim 8.62 \cdot 10^{-2} \cdot (\rho_c/\rho_a)^{-1.87}$. We also found that the application of ABW model should be limited to $15 < \rho_c/\rho_a$. In the case of $1 < \rho_c/\rho_a < 15$, an unphysical shock wave propagates into the artificial bed so that the velocity and height of the current substantially deviate from the solution satisfying the correct front condition.

Keywords: pyroclastic flows, gravity currents, shallow-water equations, numerical simulation, volcanic disaster prevention

Moment tensor representation of elliptical volume sources

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A moment tensor inversion is a powerful tool to extract source information from seismic and geodetic observations. However, widely-used moment tensor representation for volumetric sources has been limited to a few basic geometries such as a sphere, a flat crack, and a cylinder. These sources give particular diagonal component ratios: $(M_1:M_2:M_3)=(1:1:1)$ for a sphere, $(1:1:3)$ for a crack, and $(2:2:1)$ for a cylinder. When different component ratios are obtained from the inversion analysis, they are interpreted as combination of these simple geometries without considering internal pressure balance.

Although the moment tensor representation for elliptical sources was obtained 30 years ago (Davis, 1986), the solution has been rarely applied in volcanology. We consider two disadvantages of Davis (1986). The one is that the theories to relate the actual volume change to moment tensor have been proposed but not unified, which has caused some confusion. The accompanying paper (Ichihara et al., 2014, this meeting) presents a unified explanation based on the representation theorem and makes a clear link among volume change, geometry, and moment tensor. In this context, we have confirmed the applicability of Davis (1986) to the observed moment tensor.

The other disadvantage is that researchers have to search in the numerical table to find a geometry fitting to the observed moment tensor. Here we develop a facilitative tool that diagnoses the diagonal part of observed moment tensors to given the aspect ratios and the apparent compressibility. In addition, if the density and the compressibility of the internal material are given, the tool estimates mass change inside the source, which is an important parameter in volcanology.

This tool will provide a reference model satisfying pressure balance and help improving the volumetric source modeling beyond the conventional kinematic summation of simple sources.

Keywords: moment tensor, volumetric source, volcano seismology, magma

Volume source representations: a possible unified explanation based on the representation theorem

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The moment tensor inversion is a powerful tool to extract source information from seismic and geodetic observations. A moment tensor for earthquake faulting has been determined and its non-diagonal components give the seismic moment (rigidity \times slip \times fault area), which is one of essential source parameters of an earthquake. The sum of the diagonal components (the trace) of a moment tensor represents volumetric change at the source. A moment tensor determined for a volcanic source frequently has non-zero trace. However, it has been failed to uniquely relate the diagonal components to the actual volume change, which remains a critical issue in volcanic seismology (Kazahaya et al., 2011). For example, two different volume changes DV and dV have been proposed for the seismic moment of a spherical source geometry; DV comes from the moment tensor definition of a seismic fault having opening displacements whereas dV is obtained from the equivalence of resultant displacement fields due to the former moment tensor and an isotropically expanding sphere in an elastostatic equilibrium.

The difference between DV and dV has been discussed in the last decade. Muller (2001) considered an open crack of a spherical shell shape and showed that DV is the volume of the opening and dV is only the part opening outward. Aki and Richards (2002) and Richards and Kim (2005) adopted Eshelby's approach which considers virtual operations consisting of cutting, stress-free transformation, elastic straining, and welding, and concluded that the difference is due to whether the volumetric change occurs in an unconfined condition (DV) or in a confined condition (dV). Kumagai et al. (2013) reconsidered this problem and concluded that the displacement field due to a spherical source does not coincide with that due to a three-perpendicular-crack source though they both are represented by isotropic moment tensors. They also extended the insights into sources in a bimaterial medium. It is worth mentioning that the approaches of AR2002 and RK2005 give a conceptual explanation on how to adjust DV to the actual volume change dV for a sphere, but not for arbitrarily shapes. Here we address how to make such adjustments for general geometries on the basis of the representation theorem. Our imaginary operation below gives a unified explanation for the two different volumetric changes and newly proposes a method of estimating dV of the inversion results for arbitrary source geometries.

We start with the representation theorem that gives the displacement field by two terms (without a body force): a surface integration on the source region with convolution of the surface displacement and the gradient of the Green's function normal to the surface (1), and that with convolution of the surface traction and the Green's function (2). Only (1) has been considered for the seismic fault because (2) vanishes due to the balance of the traction at the contacting surfaces of a fault. On the other hand, (2) does not vanish in the case of a volumetric source, and therefore a quantitative adjustment is required to include the effect of (2) into (1). We here demonstrate that such an adjustment, is always realized by introducing an additional imaginary volumetric change, which works as 'displacement glut' in our representation of moment tensor. Our representation is found to be mathematically equivalent to the rather conceptual 'stress glut' representation proposed by Backus and Mulcahy (1976). We present a unified explanation for the existing various representations and propose a method to practically evaluate the moment tensor components from the boundary conditions of the volumetric source. The proposed representation will be useful in connecting dynamical models of volcanic processes and moment tensors.

Keywords: Moment tensor, Volumetric source, Representation theorem, Green's function, Volcano seismology, Explosion source

X-ray CT observation of delayed fragmentation of vesicular magma analog

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A laboratory experiment was performed to understand the mechanism of fragmentation of vesicular magma, which is a trigger of explosive volcanic eruption. From the observation of Shinmoe-dake 2011 eruption, the viscosity of magma is not so high as its response is in solid manner. Thus, we aim to reveal the mechanism of brittle-like fragmentation of the magma which behaves in fluid manner.

Rapid decompression test was conducted using syrup containing oxygen bubbles as a magma analogue. The decompression facility consists of a high-pressure chamber in which the specimen can be placed, and an electromagnetic valve. The pore structure of interior of the specimen was observed by X-ray microscopic tomography. The X-ray tomographic microscopy was performed at the BL20B2 beamline of the Japan Synchrotron Radiation Research Institute (JASRI, Hyogo, Japan). Initial structure of the specimen was observed by three-dimensional tomographic imaging. A digital charge-coupled device (CCD) camera was used as the detector whose imaging area is about 2048 pixels (horizontal) by 644 pixels (height) with spatial resolution of 8 $\mu\text{m}/\text{pixel}$. We took 1800 projections over 180 degrees of rotation for tomographic imaging. High-speed radiography was performed during the decompression. The framing rate of radiography is 200 frames per second.

The specimen has a semi-spherical shape whose diameter was 20 mm. The initial pressure (p_0) was 1.5 MPa, the characteristic time of decompression (t_{dec}) was 50 ms, the viscosity of syrup was about 1×10^8 Pa s, the initial averaged void fraction ϕ_0 was about 10%. We tested nine samples whose pore structure was different to each other.

The fragmentation occurred with two samples in which a 10-mm-diameter bubble was contained, while the fragmentation did not occur using the seven samples whose pore structure was relatively uniform. The onset of fragmentation is 960 ms after the decompression was started. The onset was substantially delayed not only from the characteristic decompression time but also from the relaxation time of Maxwellian viscoelastic material (viscosity/rigidity = 150 ms). Referring to the radiographic images, we found that the fracture was initiated from a chain of small gas bubbles and a notch in the vicinity of the 10-mm (large) bubble. The internal crack may grow in ductile manner when the hoop stress around the large bubble increased due to decompression. At a certain instance the stress concentration and the brittleness at crack tip may exceed the critical level, which leads to brittle failure of the crack.

The experimental result indicates that brittle-like fragmentation can occur in the non-uniform vesicular magma even if the response of magma is in fluid manner.

Model experiments on magma migration in a viscoelastic host rock : effect of viscosity

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Magma generated from partial melting ascends upwards by its buoyancy to the magma chamber, some of which erupt at the Earth's surface. Magma is considered to ascend in the form of diapir in the asthenosphere and transforms to dykes in the lithosphere (Rubin, 1995). These two end member cases have been studied in detail, but the mechanism of magma ascent in the transitional regime is still poorly known. We have been studying the ascent mechanism in the transitional regime by model experiments using a viscoelastic agar (Sumita and Ota, 2011). Here we report the results of experiments which focus on the effect of fluid viscosity on the magma migration in a viscoelastic medium.

We conducted (1) rheology measurement of agar and (2) fluid injection experiments. We inject magma (CsCl solution to which a thickener is added) using a syringe from the top of a cylindrical acrylic tank (height of 250 and 500 mm). The fluid has a volume of 1ml, density difference with the agar of 0.108 g/ml, and is injected at a constant rate of 0.1 ml/s. We vary the agar concentration in the range of 0.04-0.5 wt% and the fluid viscosity in the range of 10^{-3} - 650 Pas. As we increase the agar concentration in this range, we find that the yield stress and the rigidity of the agar increases by 3 orders of magnitude. By shearing the sample under a constant stress (creep test) we find that the agar can be approximated by a Voigt model to which a spring is connected in series. The experiments are recorded using video cameras from two sides and from the bottom of the tank.

From the injection experiments, we find that as the agar rigidity (G) decreases, the crack shape transforms from 2D (blade-like) to 3D (having a bulged head). The critical rigidity (G_c) of the 2D-3D transformation is around $G_c=10$ Pa, and this value becomes smaller when a high viscosity fluid is injected. The value of G_c is consistent with the rigidity estimated from the balance between the elasticity and buoyancy at which the strain becomes 1. The crack consists of a bulged head and a thin sheet-like tail, and the head becomes thinner and smaller as the crack elongates. When a stiff agar is used as the host, or when a high viscosity fluid is injected, the crack propagation stops within a certain distance from the injection point. We find that this propagation distance becomes shorter as we increase the agar yield stress or the viscosity of the injected fluid, and this result can be associated with the transformation of the crack shape to 2D. We fit the time-distance data to a power-law relation, and find that the exponent varies from $1/3$ to 1. An exponent of $1/3$ corresponds to the scaling obtained for a 2D crack which thins uniformly as it elongates (Taisne et al, 2011). An exponent larger than $1/3$ corresponds to the crack shape becoming more 3D. To conclude, we find that the magma viscosity not only slows its migration velocity, but also controls the crack shape, its deceleration, and the propagation distance until it stops.

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Keywords: viscoelasticity, Magma ascent, rheology, crack

Bubble growth and resorption in magma: insights from dissolved water distributions in volcanic glass

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Volcanic eruptions are driven by the growth of gas bubbles in magma, which grow and shrink as volatile species exsolve from and dissolve back into the melt in response to changes in the local environment, particularly in pressure and temperature. This movement of volatiles, particularly water, is recorded in the glass around vesicles and recent studies have used this record to interpret natural samples. Here we investigate the processes that control bubble growth and resorption in magma, by measuring the distribution of dissolved water in experimentally-vesiculated volcanic glasses. Water concentration profiles obtained using SIMS-calibrated BSEM imaging and water speciation data obtained using FTIR spectroscopy, are interpreted in the context of the known pressure and temperature history of the samples.

Samples are found to have undergone partial bubble resorption during the quench to glass at the end of experiments, as a result of increasing water solubility with decreasing temperature. Analysis of the lengthscale and timing of the resulting water concentration profiles demonstrates that the majority of resorption occurs above the glass transition. This quench resorption is associated with a reduction in bubble volumes that creates characteristic textures, such as buckled melt films between adjacent vesicles and reoriented cracks around resorption halos. Highly disequilibrium water speciation ratios within resorption halos are found to be diagnostic of quench resorption and can preserve evidence of pre-quench bubble growth

Quench resorption can increase sample water concentrations and ratios of molecular to hydroxyl water species, and reduce bubble volumes and sample porosities. Studies based on these parameters must therefore consider the potential impact of quench resorption, which is expected to be greatest for samples with high water concentrations, slow quench and low initial sample porosities. Water speciation data offer a way to investigate these impacts in unconstrained natural samples and could provide a tool for forensic interrogation of their eruptive history.

Keywords: Bubble growth, Bubble resorption, Water speciation, FTIR, SIMS, Backscatter SEM

Formulation of the 1-D magmatic flow including vesiculation kinetics

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In most of existing fluid mechanical modeling and numerical calculations, the equilibrium vesiculation has been assumed according to the solubility relation. However, in order to understand the transient behaviors such as triggering of eruptions and shifts of eruption intensity, we need to examine the effect of vesiculation kinetics on the fluid mechanical behaviors in the conduit. In this study, we formulate the governing equations describing the one dimensional fluid mechanics taken into account the vesiculation kinetics, that is, the nucleation and growth of bubbles, assuming the homogeneous flow.

As we adapt the single fluid approximation, we build upon the advantage of the formulation in which the mass conservation equation about density is converted to the pressure equation through the equation of state. We make the similar procedure to obtain pressure equation while in our case the density is a function of not only pressure but also gas phase fraction. The gas phase fraction is calculated by the vesiculation kinetics. As a result we have four equations; 1) pressure equation (mass conservation), 2) equation of motion (momentum conservation), 3) constitutive equations describing vesiculation kinetics. These partial differential equations consisting of pressure, velocity and gas phase fraction, can be solved numerically.

In order to confirm the validity and fundamental characters of the formulation, we numerically solve the shock tube problem, using modified CIP method for advective terms. In the case that the kinetic effect is negligible due to relatively large initial gas fraction in the high pressure region, we obtain the similar solution to that of single fluid, consisting of shock wave in downstream and rarefaction fan in upstream. On the other hand, in the case that kinetics of vesiculation works effectively, we have different behaviors in the high pressure region in which rarefaction front, nucleation pressure front and nucleation event propagate with different velocities, while the behavior in the low pressure region is basically same as the case without the kinetic effect. Each propagation velocity depends on nucleation pressure (liquid/gas interfacial energy), bubble growth rate, initial gas fraction in the high pressure region. The essential difference between two cases, the pressure of bubble formation becomes lower in the case with the kinetic effect than in the case without kinetic effect. If the kinetic effect is dominated, it is expected that the time to the vent from the bubble nucleation is relatively short. On the other hand, in the case that the vesiculation proceeds at equilibrium without the kinetic effect, it needs longer time to the vent, suggesting that the relative motion between gas and liquid, the bubble coalescence and degassing become dominate. Thus, this suggests that the kinetics of vesiculation may control the transition of eruption styles such as explosive and non-explosive eruptions.

Keywords: conduit flow, vesiculation kinetics, bubble nucleation, shock tube

Pre-eruptive conditions of rhyolitic magma from Kawagodaira Volcano, Izu Peninsula

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Kawagodaira volcano is a rhyolitic monogenic volcano located in the Higashi-Izu monogenic volcano field at Izu Peninsula. The volcano erupted at ca. 3060~3190. During the eruption, transition of eruption style was occurred from Plinian explosion to obsidian extrusion. Such transition often occurred in silicic magma, and depend on pre-eruptive condition and outgassing during magma ascent in the vent. However, detailed mechanism is still under study. In this study, we collect volcanic products of 3 eruption style which are plinian eruption, pyroclastic flow, and lava flow. We conducted petrography and analysis of chemical compositions of glass and phenocryst minerals for these volcanic products, and we estimated the pre-eruptive condition of temperature and pressure and water content to examine mechanism in transition of eruption style.

All samples in this study include ca. 15 vol. % of phenocrysts chiefly composed of hornblende and plagioclase with minor amount of orthopyroxene and magnetite, and include ca. 85 vol. % of groundmass composed of glass. Hornblende phenocrysts show chemical zoning with obvious gap of Al content near rim. Al content is higher in core part (ca. 1.7 atoms per formula unit (apfu)) than rim part (ca. 1.2 apfu).

The estimated pressure using hornblende Al content geobarometer shows bimodal distribution with ca. 100 MPa and 200-300 MPa in each sample. The former and the latter correspond to the crystallization pressure of rim and core part of hornblende, respectively. The estimated temperature using plagioclase-hornblende geothermometer indicate constant in all samples, and average to be 1132K. The water content is estimated to be 5wt. % by using plagioclase-melt geohydrometer, and this value coincident with H₂O solubility in rhyolitic melt at 100MPa.

The estimated P-T conditions show that magma decompressed from 200~300MPa to 100MPa without change of magma temperature. This indicates that transition of eruption style occurred less than 100MPa. The coincidence between water content and the H₂O solubility at 100MPa suggest that bubbles were already formed at this pressure. And, we think that magma stopped at 2.5~3km depth until growth of hornblende and plagioclase phenocrysts finished. During cessation of magma ascent, bubbles rised to upper part of magma chamber, forming bubble zoning in magma chamber. This zoning might relate to transition of eruption style.

Keywords: Higashi-Izu monogenetic volcanic field, rhyolitic magma, eruption style, hornblende